

Inviscid and Viscous

Fundamentals of Viscous Fluid Flow

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- inviscid viscous

viscous: sticky, greasy, adhesive,

friction, no-slip,

viscosity: relationship between stress and strain rate

(elasticity: relationship between stress and strain)

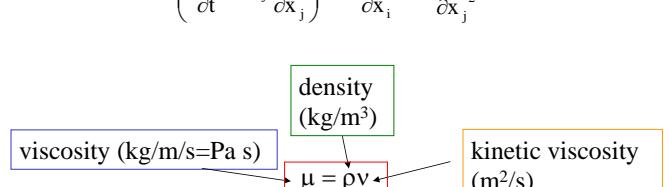
Inviscid and Viscous

- inviscid viscous
- ideal / perfect fluid real fluid
- Euler equation Navier-Stokes equation

$$\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} \quad \frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \nu \frac{\partial^2 u_i}{\partial x_j^2}$$

in tensor form

$$\rho \left(\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} \right) = - \frac{\partial p}{\partial x_i} + \mu \frac{\partial^2 u_i}{\partial x_j^2}$$



Which is more viscous, air or water?

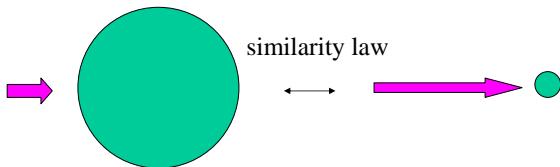
5	(m ² /s)	(kg/m ³)	μ (Pa s)
air	13.8×10^{-6} V	1.27	17.5×10^{-6} A
water	1.52×10^{-6}	1000.0	1.52×10^{-3} A

$$\frac{\partial \mathbf{u}_i}{\partial t} + \mathbf{u}_j \frac{\partial \mathbf{u}_i}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + v \frac{\partial^2 \mathbf{u}_i}{\partial x_j^2}$$

dimensionless form $\mathbf{u} = \mathbf{u}' \mathbf{U}, \quad x = x' L, \quad t = t' \frac{L}{U}, \quad p = p' \rho U^2$

$$\frac{\partial \mathbf{u}'_i}{\partial t'} + \mathbf{u}'_j \frac{\partial \mathbf{u}'_i}{\partial x'_j} = -\frac{\partial p'}{\partial x_i} + \frac{1}{Re} \frac{\partial^2 \mathbf{u}'_i}{\partial x'^2_j}$$

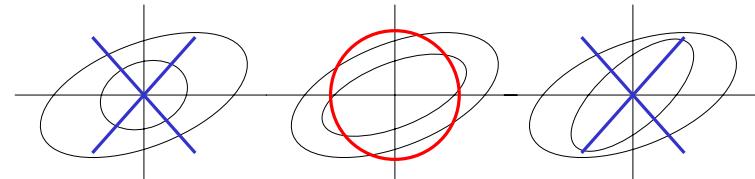
$$Re = \frac{\rho U L}{\mu} = \frac{UL}{v} \quad \text{Reynolds number}$$



Newtonian Fluid

- Newton's Law in viscosity: relationship between stress and strain rate is linear.

$$\tau_{ij} = v d_{ij} = v \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad \therefore d_{ij} = \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$



Navier-Stokes Equation

$$\begin{aligned} \frac{\partial \mathbf{u}_i}{\partial t} + \mathbf{u}_j \frac{\partial \mathbf{u}_i}{\partial x_j} &= \frac{\partial \sigma_{ij}}{\partial x_j} & \because \sigma_{ij} \equiv -p \delta_{ij} + \tau_{ij} \\ &= -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \frac{\partial \tau_{ij}}{\partial x_j} & \text{stress-divergence form} \\ &= -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \frac{\partial}{\partial x_j} \left[v \left(\frac{\partial \mathbf{u}_i}{\partial x_j} + \frac{\partial \mathbf{u}_j}{\partial x_i} \right) \right] & \therefore \frac{\partial \mathbf{u}_j}{\partial x_j} = 0 \\ &= -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + v \frac{\partial^2 \mathbf{u}_i}{\partial x_j^2} & \uparrow \end{aligned}$$

continuum equation for incompressible fluid

Irrotational Fluid

- “irrotational” means “no vortices”
- deformation of NS-equation

$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} &= -\frac{1}{\rho} \nabla p + v \nabla^2 \mathbf{u} + \mathbf{g} & \mathbf{g} \equiv (0, 0, -g) \\ \frac{\partial \mathbf{u}}{\partial t} + \nabla(\mathbf{u} \cdot \mathbf{u}) &= -\frac{1}{\rho} \nabla p + v \nabla^2 \mathbf{u} + \mathbf{g} & \because \nabla \cdot \mathbf{u} = 0 \\ \frac{\partial \mathbf{u}}{\partial t} + \frac{1}{2} \nabla(\mathbf{u} \cdot \mathbf{u}) - \mathbf{u} \times (\nabla \times \mathbf{u}) &= -\frac{1}{\rho} \nabla p + v \nabla \times (\nabla \times \mathbf{u}) + \mathbf{g} & \text{rotational form} \end{aligned}$$

$$\begin{aligned} \therefore \nabla(\mathbf{a} \cdot \mathbf{b}) &= \frac{1}{2} [\nabla(\mathbf{a} \cdot \mathbf{b}) - \nabla \times (\mathbf{b} \times \mathbf{a}) - \mathbf{b} \times (\nabla \times \mathbf{a}) - \mathbf{a} \times (\nabla \times \mathbf{b}) + \mathbf{b} \nabla \cdot \mathbf{a} + \mathbf{a} \nabla \cdot \mathbf{b}] \\ \therefore \nabla \times (\nabla \times \mathbf{a}) &= \nabla(\nabla \cdot \mathbf{a}) - \nabla \cdot \nabla \mathbf{a} \end{aligned}$$

$$\begin{aligned}\frac{\partial \mathbf{u}}{\partial t} + \nabla \left(\frac{1}{2} \mathbf{u}^2 \right) - \mathbf{u} \times \boldsymbol{\omega} &= -\frac{1}{\rho} \nabla p + \mathbf{v} \nabla \times \boldsymbol{\omega} + \mathbf{g} \\ \frac{\partial \mathbf{u}}{\partial t} + \nabla \left(\frac{1}{2} \mathbf{u}^2 + \frac{p}{\rho} - \mathbf{g} \cdot \mathbf{x} \right) &= \mathbf{u} \times \boldsymbol{\omega} + \mathbf{v} \nabla \times \boldsymbol{\omega}\end{aligned}\quad \because \boldsymbol{\omega} \equiv \nabla \times \mathbf{u}$$

if irrotational (at the same time, inviscid)

$$\begin{aligned}\frac{\partial \mathbf{u}}{\partial t} + \nabla \left(\frac{1}{2} \mathbf{u}^2 + \frac{p}{\rho} - \mathbf{g} \cdot \mathbf{x} \right) &= 0 \quad \boldsymbol{\omega} = 0 \\ \nabla \left(\frac{\partial \Phi}{\partial t} + \frac{1}{2} \mathbf{u}^2 + \frac{p}{\rho} - \mathbf{g} \cdot \mathbf{x} \right) &= 0 \quad \therefore \nabla \Phi = \mathbf{u} \\ \frac{\partial \Phi}{\partial t} + \frac{1}{2} \mathbf{u}^2 + \frac{p}{\rho} + gz &= \text{const.} \quad \Phi : \text{velocity potential}\end{aligned}$$

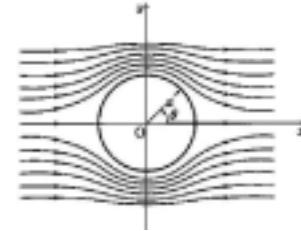
Bernoulli's theorem

$$\nabla^2 \Phi = \nabla \cdot \nabla \mathbf{u} = \nabla(\nabla \cdot \mathbf{u}) = 0$$

Laplace eq. for velocity potential = continuum eq. of incompressible fluid

What is all about Viscous Fluid?

if the flow is irrotational (and inviscid)...



$$\text{velocity potential for a cylinder in uniflow} \quad \Phi = U \left(z + \frac{a^2}{z} \right)$$

in cylindrical coordinates

$$\Phi = Ur \left(1 + \frac{a^2}{r^2} \right) \cos \theta$$

solution

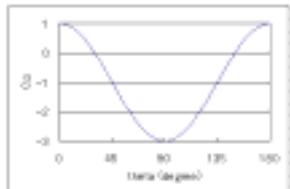
$$v_\theta = \frac{1}{r} \frac{\partial \Phi}{\partial \theta} = -U \left(1 + \frac{a^2}{r^2} \right) \sin \theta$$

on the surface of cylinder, where $r=a$

$$v_\theta = -2U \sin \theta$$

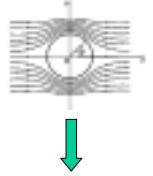
Bernoulli's theorem

$$\begin{aligned}p_0 &= \frac{1}{2} \rho (U^2 - v_\theta^2) \\ &= \frac{1}{2} \rho U^2 (1 - 4 \sin^2 \theta) \\ &= \frac{1}{2} \rho U^2 (2 \cos 2\theta - 1)\end{aligned}$$



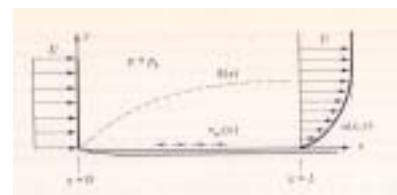
No Drag : D'Alembert's paradox

if the flow is viscous ...

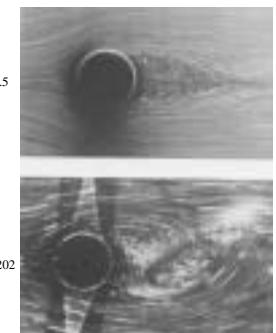


Friction

$$F = \mu \left. \frac{\partial u}{\partial y} \right|_{y=0} \quad (u = v = 0)$$



Separation



Why and how the flow separates?

NS equation

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + v \frac{\partial^2 u}{\partial y^2}$$

At the solid surface (at $y=0$)

$$(u = v = 0) \quad \frac{\partial p}{\partial x} = \mu \frac{\partial^2 u}{\partial y^2}$$

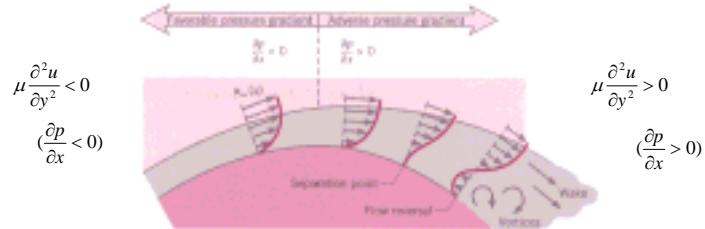


FIGURE 7.6 Velocity profile associated with separation on a circular cylinder in cross-flow.

Pressure drag

